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Introduction to stochastic dispersive dynamics

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May, 2017, RIMS, Kyoto

Introduction to stochastic dispersive dynamics ^①

Goal: Study well-posedness of stochastic dispersive PDEs.

- stochastic nonlinear Schrödinger equation (SNLS).

$$i \partial_t u = \Delta u \pm |u|^{p-1} u + \phi(\xi) \quad \pi^d \times \mathbb{R}_+$$

spacetime white noise bdd op on L^2 ($H^s: L^2 \rightarrow H^s$)

- stochastic nonlinear wave equation (SNLW)

$$\partial_t^2 u = \Delta u \pm u^p + \xi \quad \pi^d \times \mathbb{R}_+$$

Sec 1) Mild formulation & stochastic convolution (Gaussian)

$\xi(x, t) =$ space time white noise on $\pi^d \times \mathbb{R}_+$

$$\mathbb{E}[\xi(x_1, t_1) \overline{\xi(x_2, t_2)}] = \delta(x_1 - x_2) \delta(t_1 - t_2)$$

SNLS: Ito formulation

$$i du = (\Delta u \pm |u|^{p-1} u) dt + \phi dW$$

$$dW = \xi$$

$W(t) = L^2$ -cylindrical Wiener process (BM)

$$= \sum_{n \in \mathbb{Z}^d} \beta_n(t) e_n(x)$$

$\{\beta_n\} = \text{mutually indep } \mathbb{C}\text{-valued Brownian motion} \quad (2)$

$$\beta_n = \text{Re } \beta_n + i \text{Im } \beta_n$$

$\uparrow \quad \uparrow$
indep \mathbb{R} -valued BM.

(For SNLW (\mathbb{R} -valued),
we impose
 $\beta_{-n} = \overline{\beta_n}$)

For fixed $t > 0$.

$$\begin{aligned} \mathbb{E} \|W(t)\|_{H^s}^2 &= \sum_{n \in \mathbb{Z}^d} \langle n \rangle^{2s} \mathbb{E} [|\beta_n(t)|^2] \\ &\sim \sum_{n \in \mathbb{Z}^d} \langle n \rangle^{2s} < \infty \end{aligned}$$

iff $s < -\frac{d}{2}$

• Recalling the regularity of BM (in t)

$$W \in C^\alpha(\mathbb{R}_+; H^{-\frac{d}{2}-\varepsilon}(\mathbb{T}^d)), \quad \alpha < \frac{1}{2}$$

almost surely.

• Mild formulation (= Duhamel formulation) ^③

$$u(t) = S(t) u_0 \mp i \int_0^t S(t-t') |u|^{p-1} u(t') dt'$$

$$\underbrace{- i \int_0^t S(t-t') \phi dW(t')}_{= \text{Stochastic convolution}}$$

$$\Psi(t)$$

$$S(t) = e^{-it\Delta}$$

In the following, we assume that

$$\phi \in \text{HS}(L^2; H^s) \quad \text{Hilbert-Schmidt.}$$

$$\left(\begin{array}{l} \text{Recall } \phi \in \text{HS}(X; Y) \\ \text{iff } \left(\sum \| \phi e_n \|_Y^2 \right)^{1/2} < \infty, \{e_n\}, \text{ ONB of } X \end{array} \right.$$

Moreover, we assume that ϕ is homog in x .

$$\text{i.e. } \phi e_n = \hat{\phi}_n e_n$$

$$\Rightarrow \| \phi \|_{\text{HS}(L^2; H^s)} = \| \phi \|_{H^s}$$

(Banach setting: replace HS by γ -radonifying operator.)

$$\Rightarrow \mathbb{I}(t) = \sum_n e_n \hat{\Phi}_m \underbrace{\int_0^t e^{i(t-t')n^2} d\beta_n(t')}_{\text{Wiener integral}}$$

④

Wiener integral (\mathbb{R} -valued case)

Given $f \in L^2([a, b])$, consider

$$\mathbb{I}(f) = \int_a^b f dB$$

Then, $\mathbb{I}(f)$ is a mean 0 Gaussian r.v.

$$\mathbb{E}[\mathbb{I}(f)] = 0$$

with variance

$$\mathbb{E}[(\mathbb{I}(f))^2] = \|f\|_{L^2([a, b])}^2$$

Rmk: ① If f is C^1 , we can define it as a

Paley - Wiener - Zygmund integral

$$\mathbb{I}(f) = \int_a^b f dB = - \int_a^b f'(t) B(t) dt$$

(if $f(a) = f(b) = 0$)

② If B is \mathbb{C} -valued,

$$\mathbb{E}[(\mathbb{I}(f))^2] = 2 \|f\|_{L^2([a, b])}^2$$

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• Kolmogorov continuity criterion

$\{X_t\}$ with values in a metric space

Suppose

$$\mathbb{E} \left[d(X_s, X_t)^p \right] \leq C_0 |t-s|^{1+\alpha}$$

for some $p, \alpha > 0$

Then

$$P \left(\sup_{s \neq t} \frac{d(X_s, X_t)}{|t-s|^{\alpha/p - \varepsilon}} \geq 1 \right) \leq \frac{C_1}{\lambda^p}$$

$$\forall 0 < \varepsilon < \alpha/p$$

i.e. X_t is a.s. $(\frac{\alpha}{p} - \varepsilon)$ -Hölder continuous

ex: $B(t_2) - B(t_1) \sim \mathcal{N}(0, t_2 - t_1)$ $t_2 \geq t_1$

i.e. $\mathbb{E} \left[|B(t_2) - B(t_1)|^2 \right] = t_2 - t_1$

$$\Rightarrow \mathbb{E} \left[|B(t_2) - B(t_1)|^p \right] \sim |t_2 - t_1|^{\frac{p}{2}}$$

\uparrow
 $p^{p/2}$

$$\frac{\alpha}{p} = \frac{\frac{p}{2} - 1}{p} = \frac{1}{2} - \frac{1}{p} \rightarrow \frac{1}{2} \text{ as } p \rightarrow \infty$$

i.e. BM is a.s. $(\frac{1}{2} - \varepsilon)$ -Hölder.

↙ comes from hypercontractivity of OU semigroup
 • Wiener chaos estimate:

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$\{g_n\}$, seq of indep Gaussian r.v.'s

Given $k \in \mathbb{N}$, let $\{P_j\}_{j \in \mathbb{N}}$ be a seq of polynomials in $\bar{g} = \{g_n\}$ of deg at most k . Then, for $p \geq 2$,

$$\left\| \sum_{j \in \mathbb{N}} P_j(\bar{g}) \right\|_{L^p(\Omega)} \leq (p-1)^{k/2} \left\| \sum_{j \in \mathbb{N}} P_j(\bar{g}) \right\|_{L^2(\Omega)}$$

Prop 1: Let $\phi \in HS(L^2; H^s)$

Then, $\Psi \in C_t^{\frac{\alpha}{2}-} W_x^{s-\alpha, r}$ a.s. (slight loss of regularity)
 $r \leq \infty$

$r=2$: $\Psi \in C_t H_x^s$, a.s.

Pf: $t \leq T$

$$\mathbb{E} [\Psi(x, t) \overline{\Psi(y, \tau)}]$$

$$= 2 \sum_n e_n(x-y) |\hat{\phi}_n|^2 \underbrace{\int_0^t e^{i(t-t')n^2} e^{-i(\tau-t')n^2} dt'}_{}$$

$$= \int_0^t e^{i(t-\tau)n^2} dt'$$

$$= t e^{i(t-\tau)n^2}$$

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\Leftarrow Apply $\langle \nabla_x \rangle^s$ and $\langle \nabla_y \rangle^s$

$$\Rightarrow \mathbb{E} \left[\langle \nabla_x \rangle^s \Psi(x, t) \overline{\langle \nabla_y \rangle^s \Psi(y, \tau)} \right]$$

$$= 2 \sum_n e_n(x-y) \langle n \rangle^{2s} |\hat{\phi}_n|^2 t e^{i(t-\tau)n^2}$$

$$\begin{aligned} \Rightarrow \mathbb{E} \left[|\langle \nabla \rangle^s \Psi(x, t)|^p \right] &\leq p^{p/2} \mathbb{E} \left[|\langle \nabla \rangle^s \Psi(x, t)|^2 \right]^{p/2} \\ &\lesssim p^{p/2} t^{p/2} \|\phi\|_{HS(L^2; H^s)}^p \end{aligned}$$

$$\begin{aligned} \Rightarrow_{p \geq r} \left\| \|\Psi(t)\|_{W_x^{s,r}} \right\|_{L^p(\Omega)} &\leq \left\| \|\langle \nabla \rangle^s \Psi(x, t)\|_{L^p(\Omega)} \right\|_{L_x^r} \\ &\lesssim p^{1/2} t^{1/2} \|\phi\|_{HS} \end{aligned}$$

$$\Rightarrow \Psi(t) \in W_x^{s,r}, \text{ a.s.} \\ \forall r < \infty$$

Remark ① $r = \infty$: By Sobolev embedding thm,

$$\begin{aligned} \left\| \|\Psi(t)\|_{W_x^{s,\infty}} \right\|_{L^p(\Omega)} &\lesssim \left\| \|\Psi(t)\|_{W_x^{s+\varepsilon,r}} \right\|_{L^p(\Omega)} \\ &\lesssim p^{1/2} t^{1/2} \|\phi\|_{HS(L^2; H^{\underline{s+\varepsilon}})} \end{aligned}$$

In other words,

$$\Psi(t) \in W_x^{s-\varepsilon,\infty} \text{ if } \phi \in HS(L^2; H^s)$$

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$$\textcircled{2} \quad P\left(\|\Psi(t)\|_{W_x^{s,r}} > \lambda\right) \leq C e^{-c \frac{\lambda^2}{t\|\Phi\|_{HS}^2}}$$

Given $h \in \mathbb{R}$, let

$$\delta_h \Psi(x, t) = \Psi(x, t+h) - \Psi(x, t)$$

$$\Rightarrow \mathbb{E}[\delta_h \Psi(x, t) \cdot \overline{\delta_h \Psi(y, t)}]$$

$$= \mathbb{E}[\Psi(x, t+h) \overline{\Psi(y, t+h)}]$$

$$- \mathbb{E}[\Psi(x, t) \overline{\Psi(y, t+h)}]$$

$$- \mathbb{E}[\Psi(x, t+h) \overline{\Psi(y, t)}]$$

$$+ \mathbb{E}[\Psi(x, t) \overline{\Psi(y, t)}]$$

$$= 2 \sum_n e_n(x-y) |\hat{\Phi}_n|^2 \left\{ (t+h) - t e^{-i h n^2} - t e^{i h n^2} + t \right\}$$

$$= F_n(t, h) \rightarrow 0 \text{ as } h \rightarrow 0.$$

$$(\text{and } |F_n(t, h)| \leq |t| + |h|)$$

Note: $|F_n(t, h)| \leq |h| + t |1 - e^{-2ih n^2}|$

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$$\lesssim |h| + t \min(1, |h| n^2)$$

$$\lesssim |h| + t |h|^\alpha |n|^{2\alpha}, \quad \alpha \in [0, 1]$$

$$\Rightarrow \| \langle \nabla \rangle^s \delta_h \Psi(x, t) \|_{L^p(\Omega)}$$

$$\leq p^{1/2} \| \langle \nabla \rangle^s \delta_h \Psi(x, t) \|_{L^2(\Omega)}$$

$$\lesssim p^{1/2} (1+T)^{1/2} |h|^{\alpha/2} \left(\sum_n \langle n \rangle^{2(s+\alpha)} |\hat{\Phi}_n|^2 \right)^{1/2}, \quad \forall t \in [0, T]$$

$$\underline{p \geq r} \quad \| \| \delta_h \Psi(t) \|_{W_x^{s, r}} \|_{L^p(\Omega)}$$

$$\lesssim C_T p^{1/2} \| \Phi \|_{HS(L^2; H^{s+\alpha})} |h|^{\alpha/2}, \quad \forall t \in [0, T]$$

Kolmogorov conti crit. (and replace $s+\alpha$ by s)

$$\Rightarrow \Psi \in C_t^{\alpha/2 - \frac{1}{p}} W_x^{s-\alpha, r} \text{ a.s. if } \Phi \in HS(L^2; H^s)$$

$$\stackrel{p \rightarrow \infty}{\Rightarrow} \Psi \in C_t^{\frac{\alpha}{2}} W_x^{s-\alpha, r} \text{ a.s. if } \Phi \in HS(L^2; H^s)$$

• When $r=2$: It suffices to study the continuity property of $\tilde{\Psi}(t) = S(-t) \Psi(t)$

$$\Rightarrow \mathbb{E} [\tilde{\Psi}(x, t) \overline{\tilde{\Psi}(y, \tau)}] = 2 \sum_n e_n(x-y) |\hat{\Phi}_n|^2 \cdot t.$$

By repeating a similar computation,

$$\| \| \delta_h \Psi(t) \|_{H_x^s} \|_{L^p(\Omega)} \lesssim C_T p^{1/2} \| \Phi \|_{HS(L^2; H^s)} |h|^{\alpha/2}$$

Kolmogorov conti criterion
 \Rightarrow

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$$\tilde{\Psi} \in C_t^{\frac{1}{2}-} H_x^s, \text{ a.s.}$$

$$\Rightarrow \Psi \in C_t H_x^s, \text{ a.s.} \quad \begin{pmatrix} \Psi(t+h) - \Psi(t) \\ = S(t+h)(\Psi(t+h) - \Psi(t)) \\ + (S(t+h) - S(t))\Psi(t) \end{pmatrix}$$

\square

ex: $\phi = \text{Id} \Rightarrow \phi \in HS(L^2; H^s), s < -\frac{d}{2}$

$$\Rightarrow \Psi \in C_t^{\frac{\alpha}{2}-} W_x^{-\frac{d}{2}-\alpha-}, \infty$$

Sec 2 Pathwise local well-posedness

① SNLS: $s > \frac{d}{2}$

$$\Gamma u(t) = \Gamma_{u_0, \phi} u(t)$$

$$= S(t)u_0 + i \int_0^t S(t-t') |u|^{p-1} u'(t') dt' + \Psi(t)$$

$$\begin{aligned} \cdot \|\Gamma u\|_{C_T H^s} &\leq \|u_0\|_{H^s} + CT \|u\|_{C_T H^s}^p \\ &\quad + \|\Psi\|_{C_T H^s} \end{aligned}$$

$$\begin{aligned} \cdot \|\Gamma u_1 - \Gamma u_2\|_{C_T H^s} &\leq CT \left(\|u_1\|_{C_T H^s}^{p-1} + \|u_2\|_{C_T H^s}^{p-1} \right) \\ &\quad \times \|u_1 - u_2\|_{C_T H^s} \end{aligned}$$

$\Rightarrow \Gamma$ is a contraction in a ball

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$$B_R \subset C_T H^s$$

where $R_\omega = 2 (\|u_0\|_{H^s} + \|\Phi\|_{C_T H^s})$

provided

$$R_\omega^{p-1} T_\omega < 1$$

Note: $\lim_{T \rightarrow 0} \|\Phi\|_{C_T H^s} = 0$, a.s.

$$\Rightarrow \lim_{T \rightarrow 0} R_\omega^{p-1} T = 0, \text{ a.s.}$$

Thm 2: Let $\phi \in HS(L^2; H^s)$, $s > \frac{d}{2}$.

Then, SNLS is pathwise locally well-posed in H^s .

Note: local existence time $T = T_\omega$ is random.

② SNLW: $d=1$, $(u, \partial_t u) = (u_0, u_1)$ (13)

$$\cos(t|\nabla|) u_0 + \frac{\sin(t|\nabla|)}{|\nabla|} u_1$$

$$\Gamma u(t) = S(t) (u_0, u_1)$$

$$\pm \int_0^t \frac{\sin(t-t')|\nabla|}{|\nabla|} u^p(t') dt'$$

$$+ \underbrace{\int_0^t \frac{\sin(t-t')|\nabla|}{|\nabla|} dW(t')}_{= \Psi(t)}$$

$$= \Psi(t)$$

$$\Rightarrow \Psi \in C_t^{\frac{\alpha}{2}-} W_x^{\frac{1}{2}-\alpha-, \infty}$$

$s = \frac{1}{2} -$ $(u_0, u_1) \in H^s \times H^{s-1}$.

$$\|\Gamma u\|_{C_T H^s} \leq \|(u_0, u_1)\|_{H^s \times H^{s-1}}$$

$$+ T \|u^p\|_{C_T H^{s-1}} + \|\Psi\|_{C_T H^s}$$

$$\lesssim T \|u^p\|_{C_T L^{p+}}$$

$$= T \|u\|_{C_T L^{p+}}^p$$

$$\lesssim T \|u\|_{C_T H^s}^p$$

Thm 3: 1-d SNLW with additive space time white noise is pathwise locally well-posed

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in $H^s \times H^{s-1}$, $s = \frac{1}{2}$.

③ 1-d cubic SNLS in $L^2(\mathbb{T})$

$$\phi \in HS(L^2; L^2)$$

Fourier restriction norm method (Bourgain '93)

$$\|u\|_{X^{s,b}} = \left\| \langle n \rangle^s \langle \tau - n^2 \rangle^b \hat{u}(n, \tau) \right\|_{\ell_n^2 L_\tau^2}$$

$$\left(\begin{array}{l} \text{lin Schrödinger } i\partial_t u = \partial_x^2 u \\ \Rightarrow \text{FT in } x, t \quad -(\tau - n^2) \hat{u}(n, \tau) = 0. \end{array} \right.$$

• local-in-time version: $T > 0$

$$\|u\|_{X_T^{s,b}} = \inf \left\{ \|v\|_{X^{s,b}} : v|_{[0,T]} = u \right\}$$

Basic estimates: $\eta(t) = \begin{cases} 1 & \text{on } [0, 1] \\ 0 & \text{on } [-\frac{1}{2}, \frac{3}{2}]^c \end{cases}$

(i) homog lin $\|\eta(t) S(t) u_0\|_{X^{s,b}} \leq C_b \|u_0\|_{HS}$

$$\left(\mathcal{F}_{x,t} (\eta(t) S(t) u_0) = \hat{\eta}(\tau - n^2) \hat{u}_0(n) \right.$$

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(ii) nonhomog lin: $b > 1/2$, $\theta > 0$ small

$$\| \gamma(\frac{t}{T}) \int_0^t S(t-t') F(t') dt' \|_{X^{s,b}} \lesssim T^\theta \| F \|_{X^{s,b-1+\theta}}$$

(iii) L^4 -Strichartz estimate

$$\| u \|_{L^4_{x,t}(\mathbb{T} \times \mathbb{R})} \lesssim \| u \|_{X^{0,3/8}}$$

$$(iv) \| u \|_{C_T H^s} \lesssim \| u \|_{X^{s,b}_T}, \quad b > 1/2$$

(\Leftarrow 1-d Sobolev in time)

$$\begin{aligned} \Gamma u(t) &= \gamma(t) S(t) u_0 \mp i \gamma(\frac{t}{T}) \int_0^t S(t-t') |u|^2 u(t') dt' \\ &\quad + \mathbb{1}_{[0,1]} \bar{\Psi}(t) \end{aligned}$$

Take $X^{0, \frac{1}{2}-}_T$ -norm ($T \leq 1$).

$$\bullet \| \gamma(t) S(t) u_0 \|_{X^{0, \frac{1}{2}-}_T} \lesssim \| u_0 \|_{L^2}$$

continuity in time
needs to be shown separately

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$$\cdot \left\| \eta\left(\frac{t}{T}\right) \int_0^t S(t-t') |u|^2 u(t') dt' \right\|_{X_T^{0, \frac{1}{2}-}}$$

$$\leq \left\| \dots \right\|_{X_T^{0, \frac{1}{2}+}}$$

$$\leq \left\| \eta\left(\frac{t}{T}\right) \int_0^t S(t-t') |v|^2 v(t') dt' \right\|_{X^{0, \frac{1}{2}+}}$$

(for any extension v of u)

$$\lesssim T^\theta \| |v|^2 v \|_{X^{0, -\frac{1}{2}+\theta+}}$$

$$= T^\theta \sup_{\|w\|_{X^{0, \frac{1}{2}-\theta-}}=1} \left| \iint w \cdot |v|^2 v dx dt \right|$$

$$\leq \|w\|_{L_{x,t}^4} \|v\|_{L_{x,t}^4}^3$$

$$\lesssim \|w\|_{X^{0, 3/8}} \|v\|_{X^{0, 3/8}}^3$$

$$\lesssim T^\theta \|v\|_{X^{0, 3/8}}^3$$

Take inf over v .

$$\|II\|_{X_T^{0, \frac{1}{2}-}} \lesssim T^\theta \|u\|_{X_T^{0, 3/8}}^3$$

Note: $II \in C_t H^s$.

Let $b < 1/2$.

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Prop 4: $\| \mathbb{1}_{[0,1]} \Psi \|_{X^{0,b}} < \infty$, a.s.

$$\begin{aligned} \Rightarrow \| \Gamma u \|_{X_T^{0, \frac{1}{2}-}} &\lesssim \| u_0 \|_{L^2} + T^\theta \| u \|_{X_T^{0, \frac{1}{2}-}}^3 \\ &\quad + \underbrace{\| \mathbb{1}_{[0,1]} \Psi \|_{X^{0, \frac{1}{2}-}}}_{\leq C\omega < \infty \text{ a.s.}} \end{aligned}$$

Thm 5: Let $\Phi \in HS(L^2; L^2)$

Then, 1-d cubic SNLS is pathwise locally well-posed in $L^2(\mathbb{T})$.

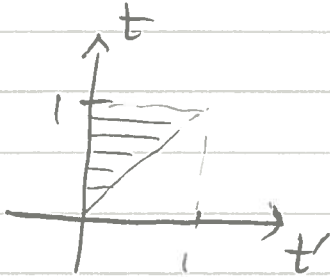
Proof of Prop 4: Suppose $\Phi \in HS(L^2; H^s)$

$$\Psi(t) = \sum_n e^{inx} \widehat{\Phi}_n \int_0^t e^{i(t-t')n^2} d\beta_n(t')$$

$$\widehat{\mathbb{1}_{[0,1]} \Psi}(m, t) = \mathbb{1}_{[0,1]} \widehat{\Phi}_m \int_0^t e^{i(t-t')m^2} d\beta_m(t')$$

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$$\widehat{1_{[0,1]} \Psi(m, \tau)} = \widehat{\Phi}_m \int_{\mathbb{R}} e^{-it(\tau - n^2)} \widehat{1_{[0,1]}(t)} \int_0^t e^{-it' m^2} d\beta_n(t') dt$$



$$= \widehat{\Phi}_m \int_0^1 e^{-it' m^2} \underbrace{\int_{t'}^1 e^{-it(\tau - n^2)} dt}_{\sim \min(1, \frac{1}{\tau - n^2})} d\beta_n(t')$$

$$\lesssim \min(1, \frac{1}{\tau - n^2})$$

$$\sim \frac{1}{\langle \tau - n^2 \rangle}$$

$$\Rightarrow \left(\mathbb{E} \| \widehat{1_{[0,1]} \Psi} \|_{X^{p, b}}^p \right)^{1/p}$$

$$\lesssim \| \langle m \rangle^s \langle \tau - n^2 \rangle^b \| \widehat{1_{[0,1]} \Psi(m, \tau)} \|_{L^p(\Omega)} \| \ell_n^2 L_\tau^2$$

$$\lesssim p^{1/2} |\widehat{\Phi}_m|^2 \frac{1}{\langle \tau - n^2 \rangle}$$

$$\lesssim p^{1/2} \left(\sum \langle m \rangle^{2s} |\widehat{\Phi}_m|^2 \int \frac{1}{\langle \tau - n^2 \rangle^{2-2b}} d\tau \right)^{1/2}$$

$$\lesssim p^{1/2} \| \Phi \|_{HS(L^2; H^s)} < \infty \text{ if } b < 1/2.$$



Da Prato - Debussche trick '03. (McKean '95
Bourgain '96)

④ Another look at 1-d cubic SNLS

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Write $u = \Psi + v$ and consider
the fixed pt problem for v .

$$v(t) = g(t) S(t) u_0 \mp i g\left(\frac{t}{T}\right) \int_0^t S(t-t') |v + \Psi|^2 (v + \Psi) dt'$$

\Rightarrow contraction in $X_T^{0, \frac{1}{2}+}$ (a bit smoother)

only change:

$$\|v + \Psi\|_{L_{x,T}^4} \leq \|v\|_{L_{x,T}^4} + \|\Psi\|_{L_{x,T}^4}$$

$< \infty$, a.s.

\Leftarrow No need to take the $X^{s,b}$ -norm of Ψ .

Remark: ϕ can be random / time dependent, etc.

Sec 3) On global well-posedness

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① 1-d cubic SNLS in $L^2(\mathbb{T})$

• deterministic NLS

LWP in $L^2(\mathbb{T})$ & L^2 -conservation

\Rightarrow By iterating LWP, we obtain GWP in $L^2(\mathbb{T})$

Idea: Control $M(u) = \sum_n |\hat{u}_n|^2 = \sum_n (p_n^2 + q_n^2)$

$$p_n = \operatorname{Re} \hat{u}_n, \quad q_n = \operatorname{Im} \hat{u}_n$$

$$d\hat{u}_n = \left(i n^2 \hat{u}_n - i \mathcal{F}_x(|u|^2 u)(n) \right) dt \\ - i \hat{\Phi}_n d\beta_n$$

$$\Rightarrow dp_n = \left(-n^2 q_n + \operatorname{Im} \mathcal{F}_x(|u|^2 u)(n) \right) dt$$

$$+ \underbrace{\operatorname{Im}(\hat{\Phi}_n d\beta_n)} \\ = \operatorname{Im} \hat{\Phi}_n \cdot d(\operatorname{Re} \beta_n) \\ + \operatorname{Re} \hat{\Phi}_n d(\operatorname{Im} \beta_n)$$

$$dq_n = \left(n^2 p_n - \operatorname{Re} \mathcal{F}_x(|u|^2 u)(n) \right) dt$$

$$- \underbrace{\operatorname{Re}(\hat{\Phi}_n d\beta_n)} \\ = -\operatorname{Re} \hat{\Phi}_n d(\operatorname{Re} \beta_n) + \operatorname{Im} \hat{\Phi}_n d(\operatorname{Im} \beta_n)$$

Ito's lemma: $dX = f dt + g dB$

Consider $F(X)$

$$\text{Then, } dF = \partial_x F dX + \frac{1}{2} \partial_x^2 F (dX)^2$$

$$= \partial_x F (f dt + g dB)$$

$$+ \frac{1}{2} \partial_x^2 F \cdot g^2 dt$$

$$(dt)^2 = 0$$

$$dt dB = dB dt = 0$$

$$(dB)^2 = dt$$

$$\Rightarrow dM \stackrel{\text{Ito}}{=} 2 \sum (p_n dp_n + q_n dq_n)$$

$$+ \sum \left((dp_n)^2 + (dq_n)^2 \right)$$

$$= \sum \left\{ \left(\text{Im} \hat{\Phi}_n d(\text{Re} \beta_n) + \text{Re} \hat{\Phi}_n d(\text{Im} \beta_n) \right)^2 + \left(\text{Re} \hat{\Phi}_n d(\text{Re} \beta_n) + \text{Im} \hat{\Phi}_n d(\text{Im} \beta_n) \right)^2 \right\}$$

$$= 2 \sum |\hat{\Phi}_n|^2 dt$$

$$= 2 \|\Phi\|_{\text{HS}(L^2; L^2)}^2 dt$$

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$$\begin{aligned}
& 2 \sum (p_n dp_n + q_n dq_n) \\
&= 2 \sum \left\{ p_n (-n^2 q_n + \operatorname{Im} \mathcal{F}_x(|u|^2 u)(n)) dt \right. \\
&\quad \left. + q_n (n^2 p_n - \operatorname{Re} \mathcal{F}_x(|u|^2 u)(n)) dt \right\} \overset{L^2 \text{-conservation}}{\downarrow} = 0 \\
&+ 2 \sum_n \left(p_n \operatorname{Im}(\hat{\Phi}_n d\beta_n) - q_n \operatorname{Re}(\hat{\Phi}_n d\beta_n) \right)
\end{aligned}$$

Burkholder - Davis - Gundy inequality:

X , (local) martingale

$1 \leq p < \infty$,

$$\mathbb{E} \left[\sup_{t \in [0, T]} |X_t|^p \right] \sim \mathbb{E} \left[\langle X \rangle_{[0, T]}^{p/2} \right]$$

quadratic variation

For Ito processes $dX = f dt + g dB$

$$\langle X \rangle_{[0, t]} = \int_0^t g^2 dt'$$

(22)

One of the terms is given by

$$2 \sum_n p_n \operatorname{Im} \hat{\Phi}_n d(\operatorname{Re} \beta_n)$$

$$\Rightarrow \text{Take } \mathbb{E} \left[\sup_{t \in [0, T]} \int_0^t \cdot dt' \right]$$

and apply B-D-G ineq,

$$\lesssim \mathbb{E} \left[\left(\int_0^T \sum_n p_n^2 |\hat{\Phi}_n|^2 dt \right)^{1/2} \right]$$

$$\leq \|u\|_{L^2}^2 \|\Phi\|_{L^2}^2 \quad \left(\begin{array}{l} \sum a_n b_n \\ \leq \sum a_n \sum b_n \end{array} \right)$$

$$\leq \mathbb{E} \left[\left(\sup_{t \in [0, T]} M(u)(t) \right)^{1/2} \cdot T^{1/2} \|\Phi\|_{L^2} \right]$$

$$\leq \underbrace{\varepsilon \mathbb{E} \left[\sup_{t \in [0, T]} M(u)(t) \right]}_{\text{hide in the left-hand side}} + \frac{1}{\varepsilon} T \mathbb{E} \left[\|\Phi\|_{L^2}^2 \right]$$

Putting together, we obtain

$$\mathbb{E} \left[\sup_{t \in [0, T]} \|u(t)\|_{L^2}^2 \right] \lesssim \|u_0\|_{L^2}^2 + C(T, \|\Phi\|_{\text{HS}(L^2; L^2)})$$

Thm 6: 1-d cubic SNLS is globally well-posed in $L^2(\Pi)$

(Similar GWP in H^1 holds for energy-subcritical SNLS in the defocusing case.)

(23)

Rmk: In order to justify the application of Itô's lemma, one should first consider the finite dim'l approximation

$$\begin{cases} i d u_N = (\Delta u_N \pm P_{\leq N} (|u_N|^2 u_N)) dt + P_{\leq N} (\phi dW) \\ u_N|_{t=0} = P_{\leq N} u_0 \end{cases}$$

for which (finite dim'l) Itô's lemma holds.

\Rightarrow obtain an a priori bound (indep of N)

Then, proceed with an approximation argument

$u_N \rightarrow u$ (a variant of the LWP argument)

and obtain the same a priori bd for a soln u to SNLS.

• For the infinite dim'l version of Itô's lemma, see Da Prato-Zabczyk (2nd ed)

② 1-d cubic ^{defocusing} SNLW with additive space time white noise (24)

$$\begin{cases} \partial_t^2 u = \partial_x^2 u - u^3 + \mathfrak{Z} \\ (u, \partial_t u)|_{t=0} = (u_0, u_1) \in \mathcal{H}'(\mathbb{T}) \end{cases}$$

With $\Psi(t) = \int_0^t \frac{\operatorname{Im}((t-t')|\nabla|)}{|\nabla|} dW(t') \in C_t W_x^{\frac{1}{2}-, \infty}$

Write $u = v + \Psi$. Then, the residual part v satisfies

$$\begin{cases} \partial_t^2 v = \partial_x^2 v - (v + \Psi)^3 \\ (v, \partial_t v)|_{t=0} = (u_0, u_1) \end{cases}$$

By the local argument (exercise!!), we have

$$v \in C_{T_\omega} H'.$$

and

$$\lim_{t \rightarrow T_{\max}(\omega)} \|v(t)\|_{H'} = \infty \quad \text{if } T_{\max}(\omega) < \infty$$

Idea: Control $E(v) = \frac{1}{2} \int (\partial_x v)^2 + \frac{1}{2} \int (\partial_t v)^2 + \frac{1}{4} \int v^4$

(25)

$$\partial_t E(v) = \int (\partial_t v) \left(\underbrace{\partial_t^2 v - \partial_x^2 v + v^3}_{= - (v + \Phi)^3} \right) dx$$

$$= - (v + \Phi)^3$$

$$\underbrace{}$$

$$= - v^2 \Phi - v \Phi^2 - \Phi^3$$

$$= O(|v^2 \Phi| + |\Phi|^3)$$

$$\stackrel{C-S}{\lesssim} E(v)^{1/2} \left(\|\Phi(t)\|_{L_x^\infty} E(v)^{1/2} + \|\Phi(t)\|_{L_x^6}^3 \right)$$

$$\lesssim \|\Phi\|_{L_T^\infty L_x^6}^6 + \left(1 + \|\Phi\|_{L_{T,x}^\infty}^2 \right) E(v)$$

\Rightarrow By Gronwall, $E(v) \leq C(T, \Phi^w) < \infty$

Thm 7: 1-d cubic SNLW with additive space time white noise is globally well-posed in $H'(\mathbb{T})$

- One can handle the quintic nonlinearity in a similar manner but needs an integration-by-parts trick (Oh-Pocovnicu '16 for 3-d quintic NLW)

Stochastic nonlinear beam equation (SNLB) (26)

$$\partial_t^2 u = \Delta^2 u - u^p + \Xi \quad \text{on } \mathbb{T}^d$$

$$\Xi(t) = \int_0^t \frac{\sin(t-s)}{\Delta} dW(s) \in C_t W_x^{2-\frac{d}{2}-, \infty}$$

Thm 8 (Moscato-Pocovnicu-Tolomeo-Wang '17) $d=3$

Let $1 < p < 11/3$. Then, SNLB is globally well-posed in $H^2 \times L^2$.

Sec 4 Remarks & comments

① We can also consider SPDEs with a multiplicative noise.

$$(SNLS_m) \quad i\partial_t u = \Delta u \pm |u|^{p-1} u + u \phi d\Xi$$

↑
Ho or Stratonovich product

Key lemma: $p \geq 2$

$$\mathbb{E} \left[\left\| \int_0^t S(t-t') u \phi dW \right\|_{X_T^{s, \frac{1}{2}-}}^p \right] \lesssim C_{T, \phi} \mathbb{E} \left(\|u\|_{X_T^{s, \frac{1}{2}-}}^p \right)$$

\Leftarrow B-D-G inequality

(27)

As a consequence, we construct solns in

$$L^2(\Omega; X_T^{s, \frac{1}{2}-} \cap C_T H^s)$$

i.e. NOT pathwise!!

(\Leftarrow uses the "truncation" method
to handle a nonlinearity.

Q1: pathwise well-posedness of SNLS with
a multiplicative noise.

Q2: pathwise well-posedness of SNLS with
difficult a multiplicative space-time white noise

$$\begin{array}{cc} u & dW \\ \cap & \cap \\ u \in C_t^{\frac{1}{2}-} C_x^{\frac{1}{2}-} & C_t^{-\frac{1}{2}-} C_x^{-\frac{1}{2}-} \end{array}$$

$(\frac{1}{2}-) + (-\frac{1}{2}-) < 0 \Rightarrow$ cannot make sense
of the product.

② Singular stochastic dispersive PDEs

28

2.② 2-d SNLW with additive space-time white noise

$$\partial_t^2 u = \Delta u \pm u^k + \xi, \quad x \in \mathbb{T}^2.$$

Write $u = v + \Psi$ where

$$\Psi = \int_0^t \frac{\sin(t-t')|\nabla|}{|n|} dW(t')$$

$$\Rightarrow \text{formally, } \partial_t^2 v = \Delta v \pm (v + \Psi)^k$$

BUT in 2-d, we have

$$\Psi \in \underline{C_t H_x^{-\varepsilon}} \setminus C_t L_x^2$$

$\Rightarrow \Psi^k$ does not make sense and we need to renormalize the nonlinearity.

• Truncated stochastic convolution

$$\Psi_N = P_{\leq N} \Psi$$


$$= \sum_{|n| \leq N} e_n(x) \int_0^t \frac{\sin(t-t')|n|}{|n|} d\beta_n(\tau)$$

For fixed $(x, t) \in \mathbb{T}^2 \times \mathbb{R}_+$,

$$\mathbb{I}_N(x, t) = \text{mean } 0 \text{ Gaussian r.v. with}$$

variance $\sigma_N(t) = \mathbb{E}[\mathcal{I}_N^2(x, t)]$

$$n + \log N \rightarrow 0 \quad \text{as } N \rightarrow \infty$$

- For $U_N = V_N + \Psi_N$,


we have

$$u_N^k = \sum_{l=0}^k \binom{k}{l} \Psi_N^l v_N^{k-l}$$

Wick ordering replace \mathbb{I}_N^l by

$$: \Psi_N^l(x, t) : = H_l(\Psi_N(x, t); \sigma_N(t)) \text{ ptwise}$$

Hermite poly of deg l defined by

$$e^{tx - \frac{1}{2}\sigma t^2} = \sum_{l=0}^{\infty} \frac{t^l}{l!} H_l(x; \sigma)$$

Rmk: Wick ordering = orthogonal projection onto the Wiener chaoses of deg l

By repeating a computation analogous to Prop 1 on page (6), we have

Prop 9: $\{\Psi_N^l\}$ is a Cauchy seq in $L^p(\Omega; C([0, T]; W^{-\varepsilon, \infty}(\mathbb{T}^2)))$

$$\Rightarrow \Psi^l \stackrel{\text{def}}{=} \lim_{N \rightarrow \infty} \Psi_N^l$$

$$\text{By } H^k(x+y) = \sum_{l=0}^k \binom{k}{l} H^l(x) y^{k-l},$$

we have

$$\begin{aligned} \Psi_N^k(x, t) &\stackrel{\text{def}}{=} H^k(U_N(x, t); \sigma_N(t)) \\ &= \sum_{l=0}^k \binom{k}{l} H^l(\Psi_N(x, t); \sigma_N(t)) (v_N(x, t))^{k-l} \end{aligned}$$

$$\downarrow N \rightarrow \infty$$

$$\Psi^k(x, t) \stackrel{\text{def}}{=} \sum_{l=0}^k \binom{k}{l} \Psi^l(x, t) (v(x, t))^{k-l}$$

$$=: F_{\Psi}(v(t))$$

$$\text{for } u = v + \Psi$$

\nwarrow smoother

(NOT defined for all functions!!)

\Leftarrow Time-dependent renormalization

⇒ Consider the fixed pt problem for v

(31)

$$v(t) = S(t)(\phi_0, \phi_1) \mp \int_0^t \frac{\sin(t-t')|\nabla|}{|\nabla|} F_{\pm}(v(t')) dt'$$

Thm 10: (Gubinelli-Koch-Oh '17)

Wick ordered SNLW

$$(WSNLW) \quad \partial_t^2 u - \Delta u \pm :u^k: = \xi$$

is pathwise locally well-posed in $H^s(\mathbb{T}^2)$

$$(i) \quad k \geq 4; \quad s \geq s_{crit}$$

$$(ii) \quad k = 2, 3; \quad s > s_{crit},$$

where

$$s_{crit} = \max(s_{scaling}, s_{conf}, 0)$$

$$\swarrow \text{scaling critical regularity} = \frac{d}{2} - \frac{2}{k-1}$$

$$\searrow \text{conformal critical regularity} = \frac{d+1}{4} - \frac{1}{k-1}$$

$$u = \Psi + v$$

$$\in \Psi + C([0, T]; \underline{H^\sigma(\mathbb{T}^2)})$$

$$\subset C([0, T]; H^{-\varepsilon}(\mathbb{T}^2)), \quad \forall \varepsilon > 0$$

$$(\underline{\sigma = \min(s, 1 - \varepsilon)}.)$$

Idea: Solve the fixed pt problem for v by

① Strichartz estimates,

② Prop 9 ($\Psi^l: \in C([0, T]; \bar{W}^{\varepsilon, \infty}(\mathbb{T}^2))$)

③ product estimate: $\frac{1}{p} + \frac{1}{q} = \frac{1}{r} + \frac{s}{d}$

$$\|\langle \nabla \rangle^{-\varepsilon} (fg)\|_{L^r} \lesssim \|\langle \nabla \rangle^{-\varepsilon} f\|_{L^p} \|\langle \nabla \rangle^{\varepsilon} g\|_{L^q}$$

Global well-posedness?

defocusing cubic WSNLW.

$$\partial_t^2 v - \Delta v + v^3 = -3v^2 \Psi - 3v : \Psi^2 : - : \Psi^3 :$$

Two difficulties:

① $v(t) \notin H^1$

\Leftarrow can not use the energy:

$$E(v) = \frac{1}{2} \int (\partial_x v)^2 + \frac{1}{2} \int |\nabla v|^2 + \frac{1}{4} \int v^4$$

We need to smooth v .

\Rightarrow I-method in the stochastic setting.

② Even if v were in H^1 , v does not satisfy (deterministic) NLW, i.e. $E(v)$ is not conserved.

\Rightarrow Need a Gronwall-type argument as in pp. 24-25.

Thm 11: (Gubinelli-Koch-Oh-Tolomeo '17)

The defocusing cubic WSNLW is globally well-posed

2.6 Open questions (difficult)

(34)

- 1-d SNLS with space-time white noise forcing

$$i \partial_t u = \partial_x^2 u \pm |u|^2 u + \xi$$



need a renormalization

- 3-d SNLW with space-time white noise forcing

$$\xi \in C_t^{-\frac{1}{2}-} H_x^{-\frac{3}{2}-}$$

$$\Rightarrow \Psi \in C_t H_x^{-\frac{1}{2}-}$$